

# Phase Diagram of Shastry-Sutherland Compound $\text{SrCu}_2(\text{BO}_3)_2$ at High-Pressure and High-Field

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INTERNATIONAL WORKSHOP ON QUANTUM MAGNETS IN EXTREME CONDITIONS

*Please join via [Zoom](#)*

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# Acknowledgments

**Frederic Mila**



**David Graf**



**Cristian Batista**



**Philippe Corboz**



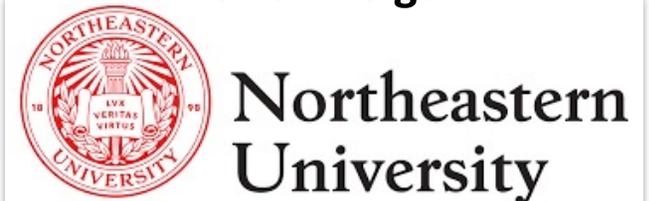
**Marcelo Jaime**



**Thomas Rosenbaum**



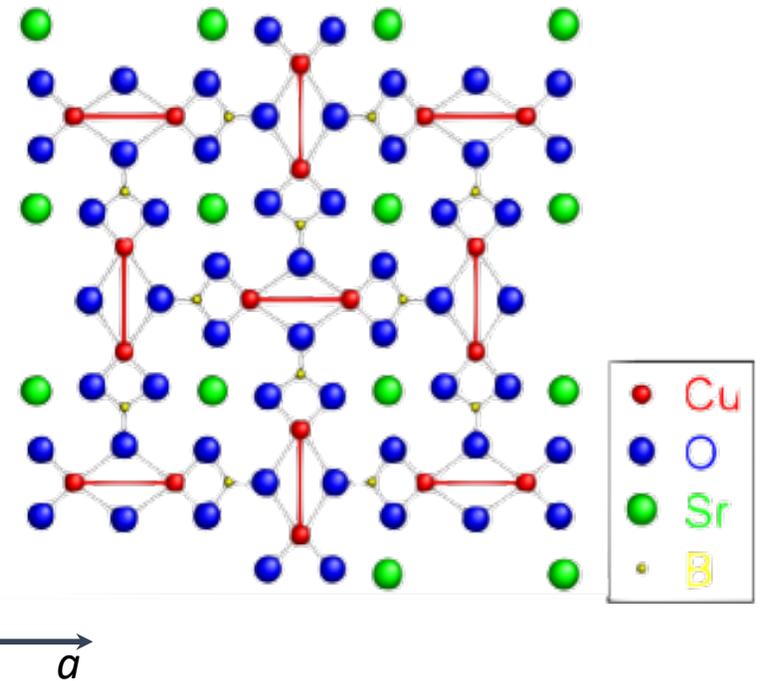
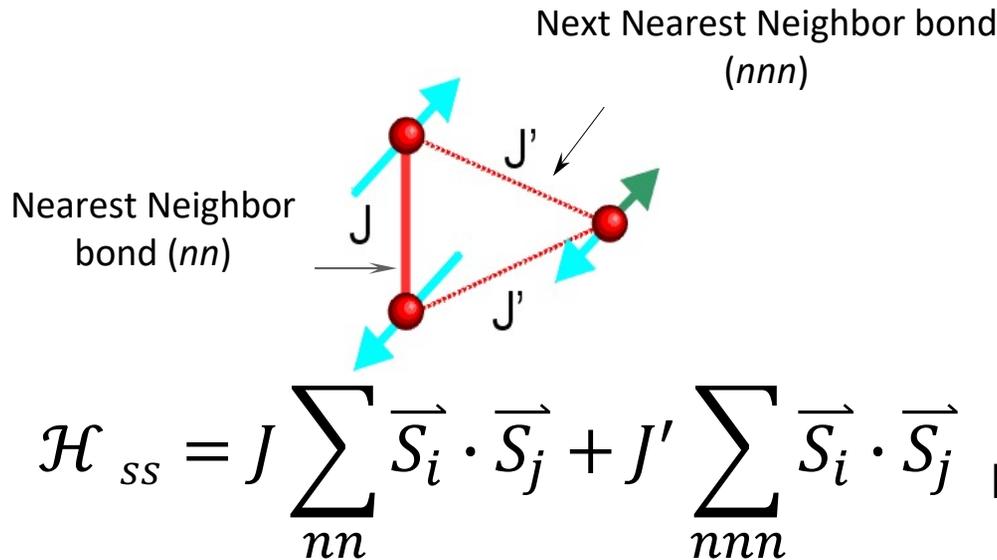
**Adrian Feiguin**





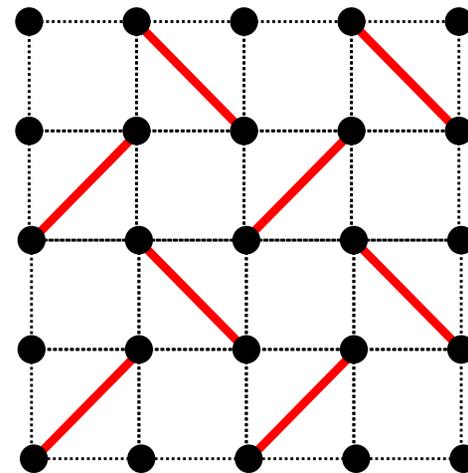
**My Team @ Duke (pre-COVID!)**

# SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub>: Realization of Shastry-Sutherland Model



- $a - b$  plane topologically equivalent to Shastry-Sutherland lattice
- $x \equiv \frac{J'}{J} \sim 0.63$
- Weak inter-plane coupling
- Isolated singlet dimer ground state

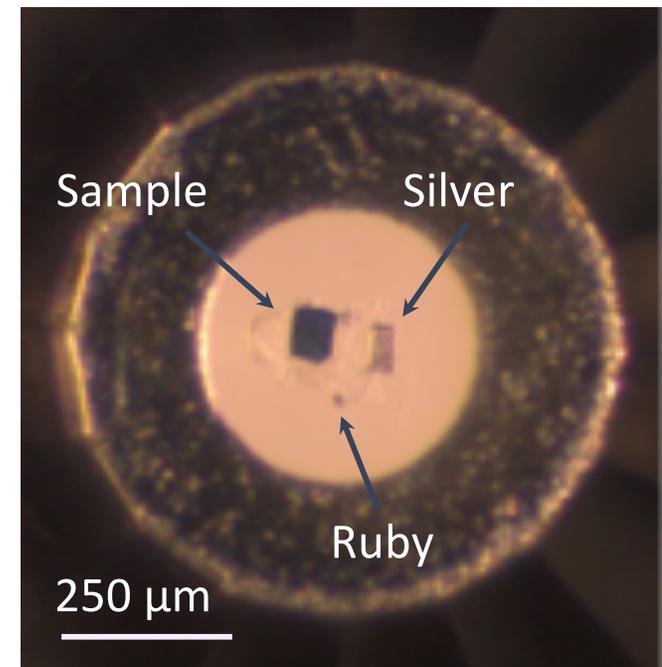
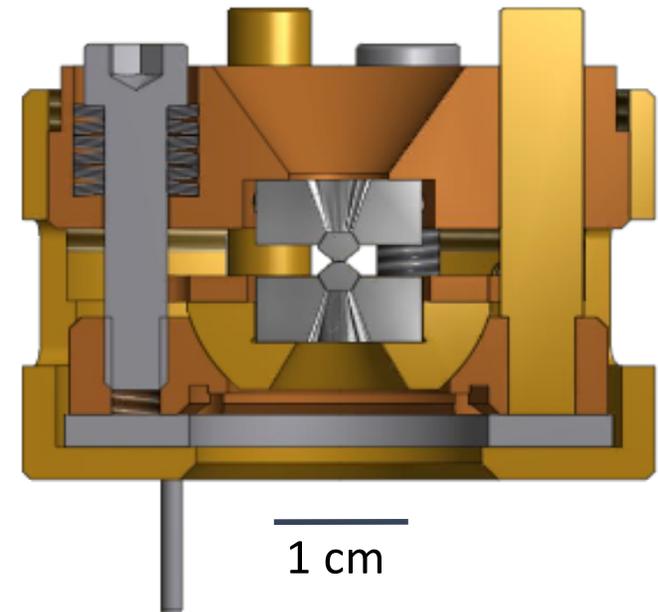
→ Use pressure to change bond distances  
vary  $x$ , tune ground state



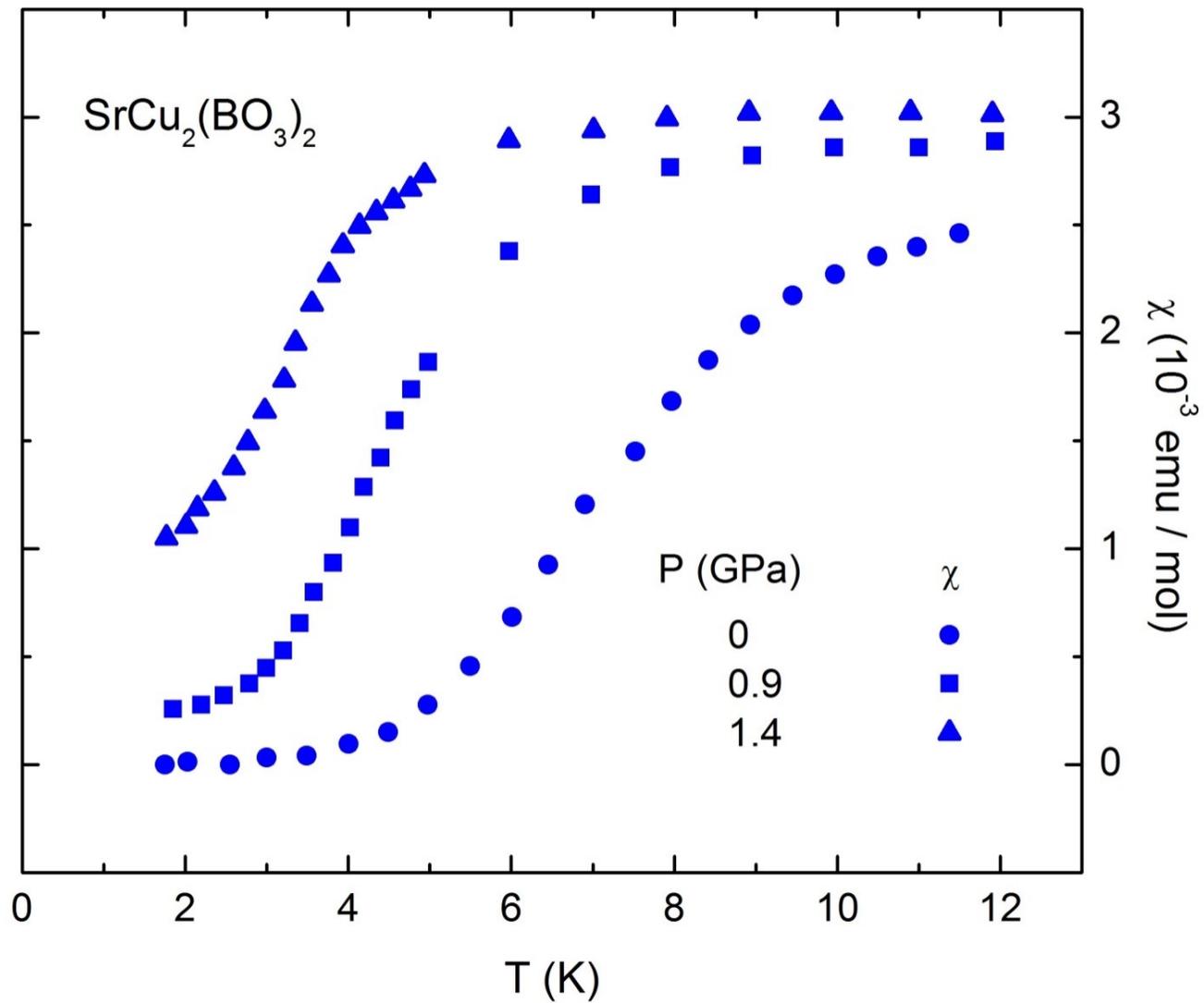
# Tuning knob 1: Pressure

## Ground states at high pressure

- ❖ Vary coupling ratio by applying hydrostatic pressure to compress sample
- ❖ Experimental range ~1-7 GPa: Diamond anvil cell
  - Single-crystal sample confined between faces of two diamonds, metallic gasket
  - Methanol/Ethanol pressure medium for hydrostaticity
  - Ruby and silver manometers
  - In situ pressure tuning via helium gas membrane
- ❖ X-ray measurements: Transmission geometry, synchrotron source

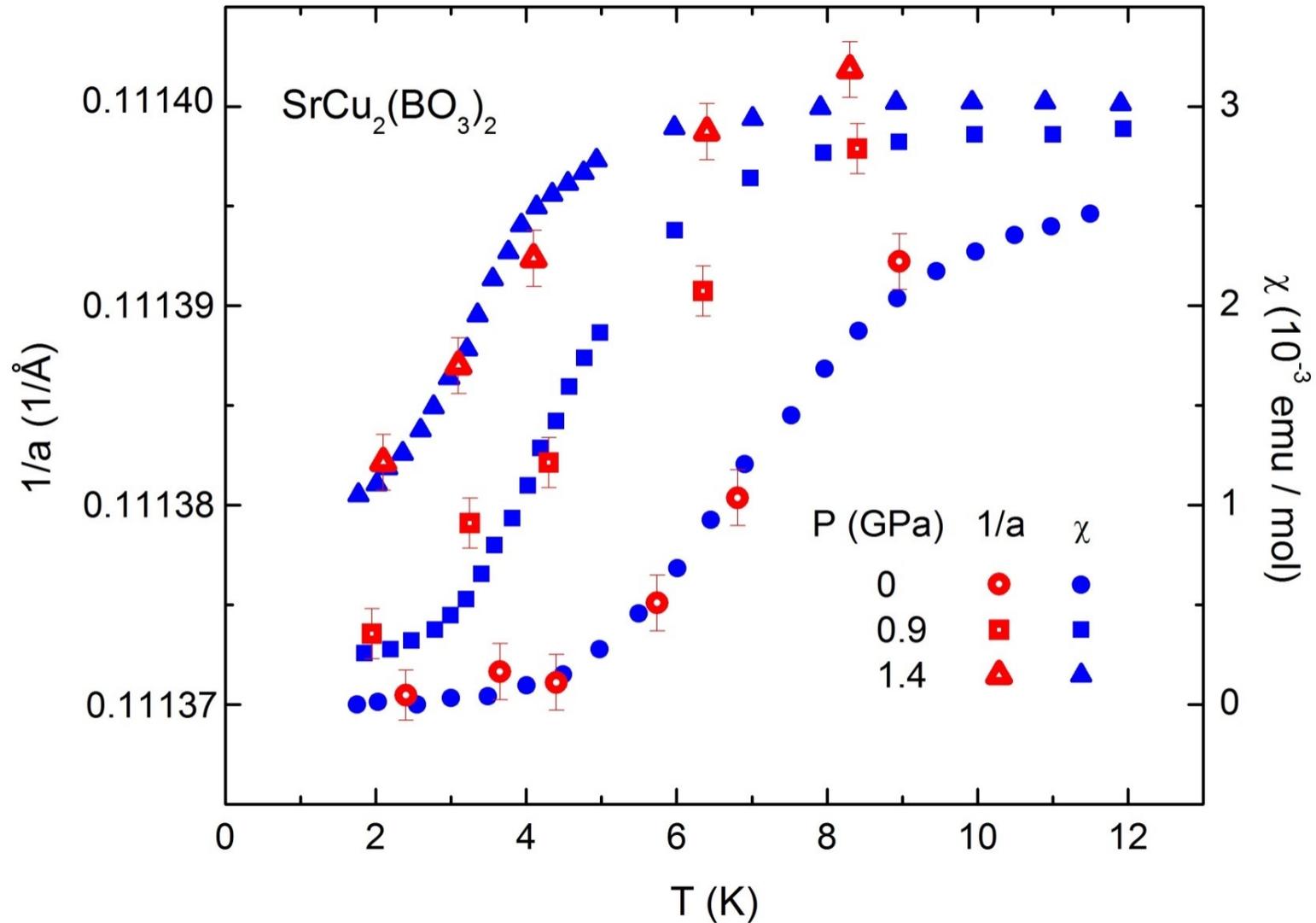


# Pressure Tuning of Energy Gap in $\text{SrCu}_2(\text{BO}_3)_2$



Waki *et al.*, *J. Phys. Soc. Jpn.*, 76, 073710 (2007).

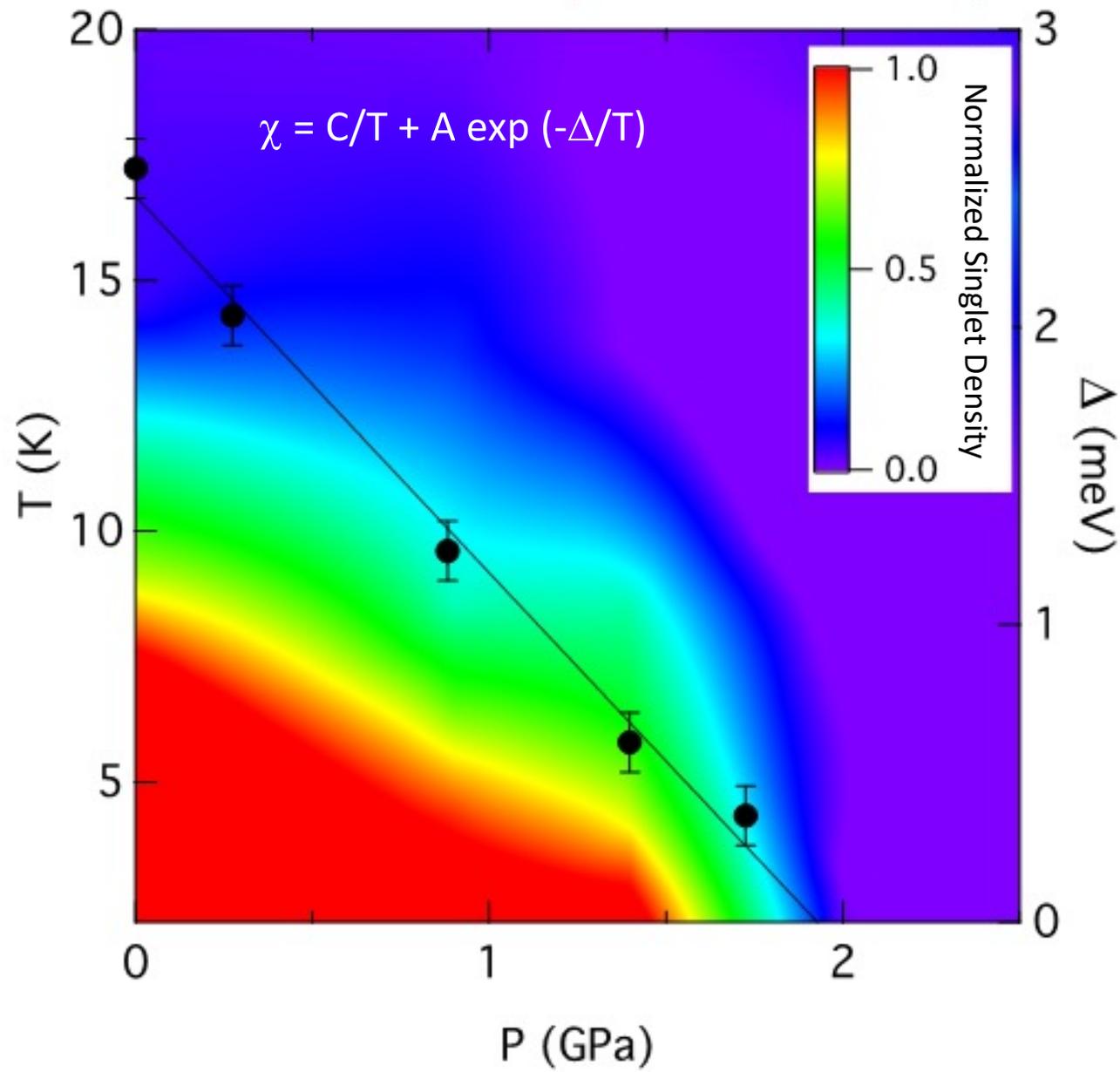
# Strong Spin-Phonon Coupling in $\text{SrCu}_2(\text{BO}_3)_2$



Waki *et al.*, *J. Phys. Soc. Jpn.*, 76, 073710 (2007).

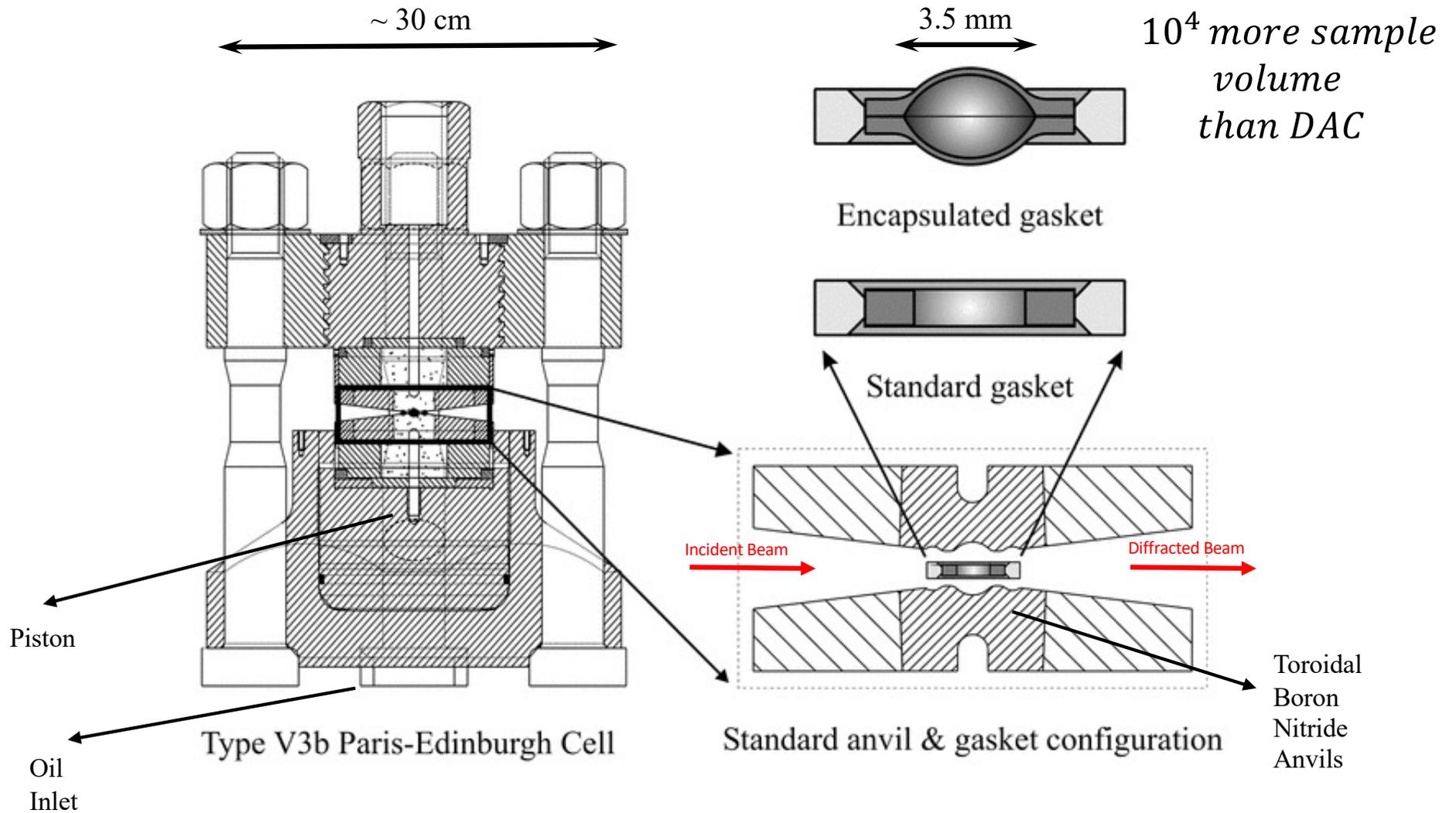
Haravifard *et al.*, *Proc. Natl. Acad. Sci. USA* 109, 2286 (2012).

# Singlet Density & First Phase Transition

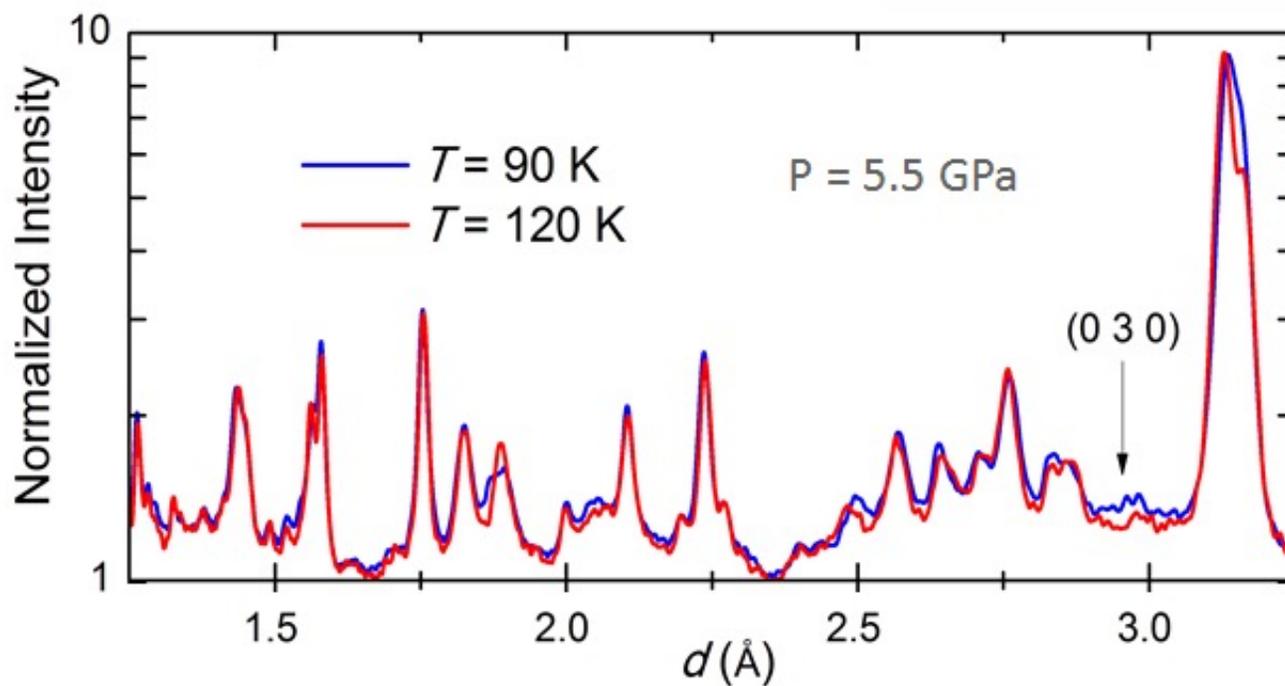
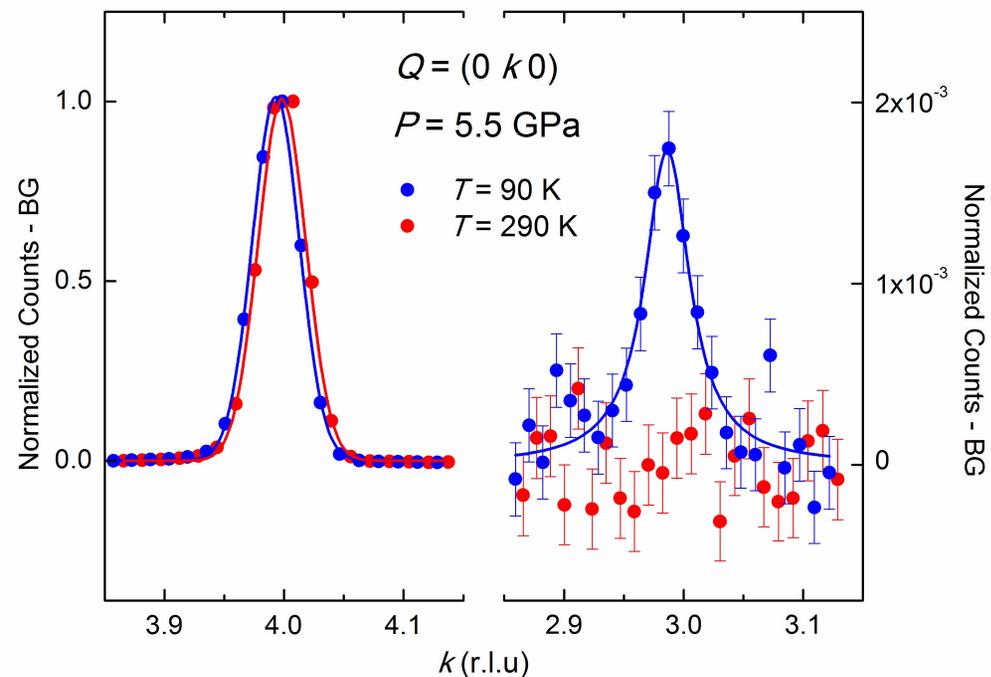
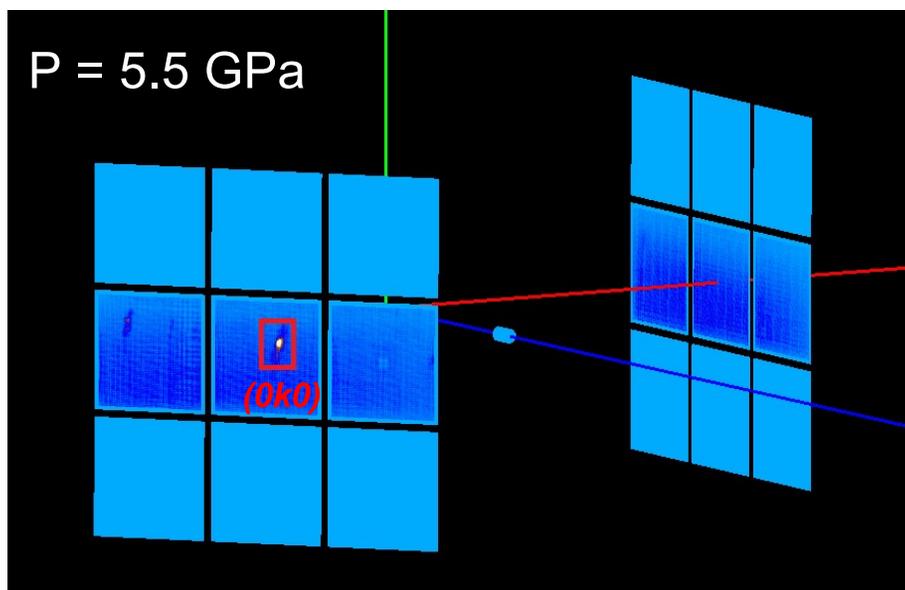


Haravifard *et al.*, *Proc. Natl. Acad. Sci. USA* 109, 2286 (2012).

# High Pressure Neutron Diffraction Measurements Paris-Edinburgh Cell

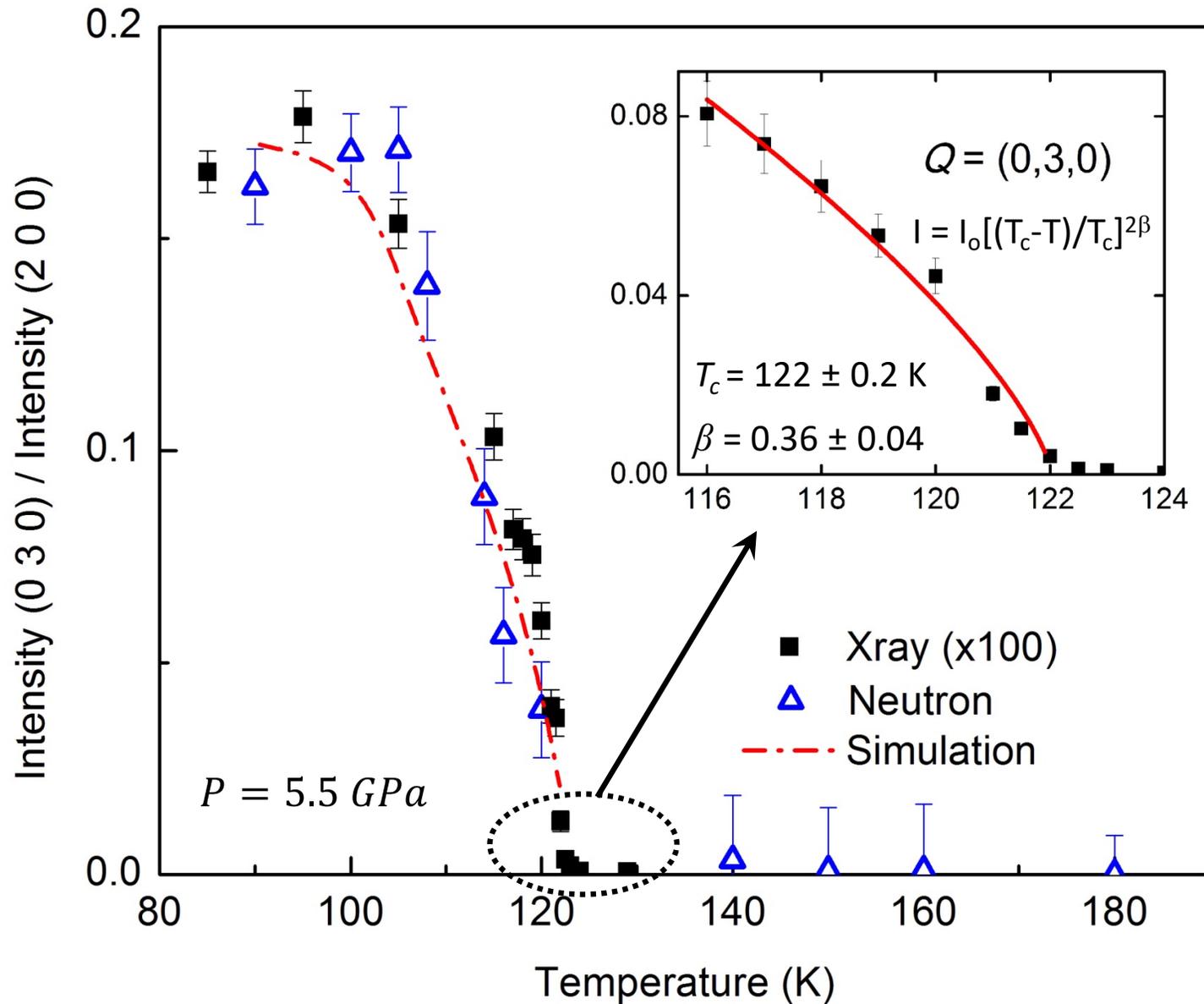


# High Pressure Neutron Diffraction Measurements Conducted at SNAP



Haravifard *et al.*,  
*Proc. Natl. Acad. Sci.*  
111, 14372 (2014).

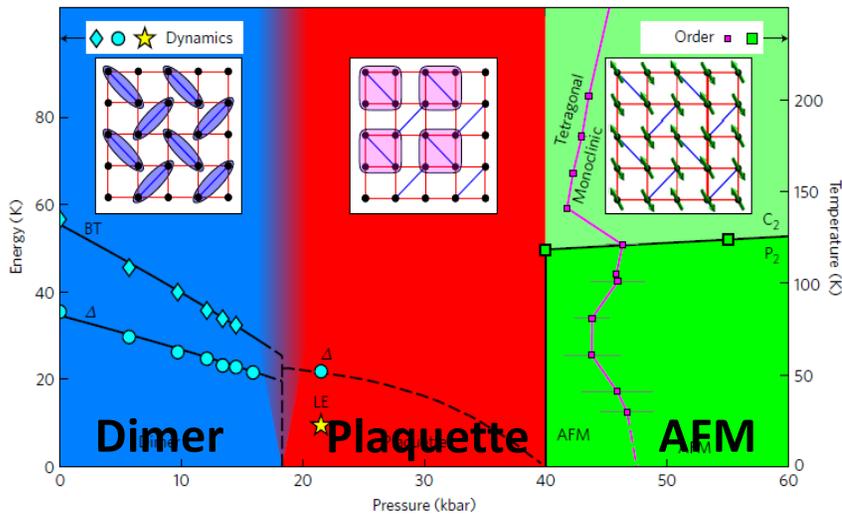
### 3-Dimensional Nature Shown by Forbidden Peaks and Critical Behavior



# Tuning knob 1: Pressure

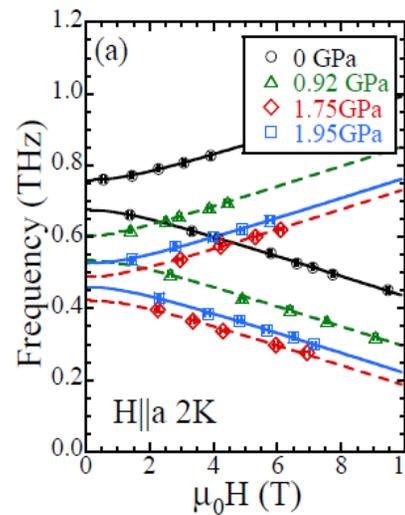
## Ground states at high pressure

### Inelastic neutron scattering



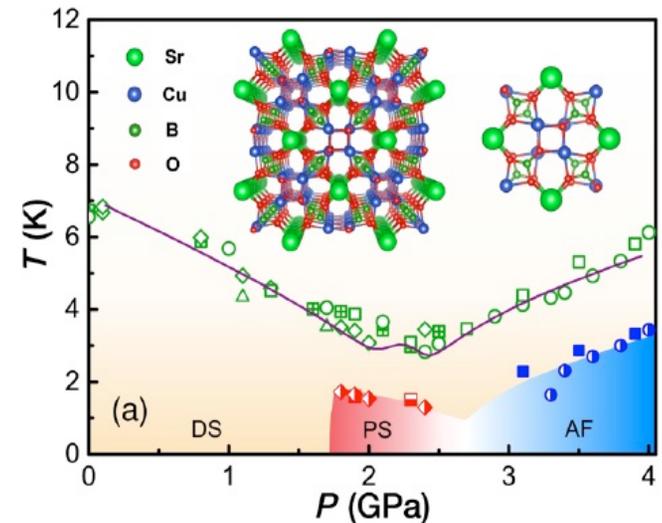
M. Zayed, *et al.*, *Nat. Phys.* **13**, 962 (2017)

### Terahertz ESR



T. Sakurai, *et al.*, *J. Phys. Soc. Jpn.* **87**, 033701 (2018)

### Specific heat



J. Guo, *et al.*, *PRL*. **124**, 206602 (2020)

Extreme environments needed

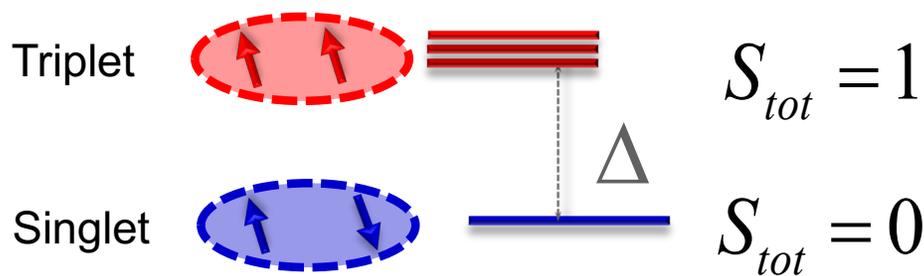


$T < \sim 2$  K  
 $P \sim 1.8$  GPa: dimer  $\rightarrow$  plaquette  
 $P \sim 2.5$  GPa: plaquette  $\rightarrow$  AFM

# Tuning knob 2: Magnetic field

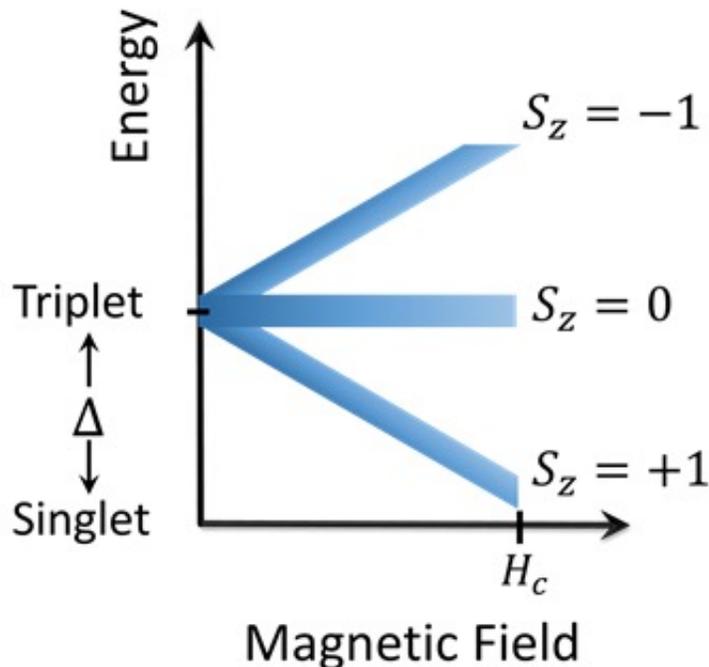
## Magnetization plateaus

Analogy can be portrayed between a spin  $\frac{1}{2}$  dimer system in an external magnetic field and a lattice gas of hard-core bosons.



$$\mathcal{H} = J \sum_{nn} \vec{S}_i \cdot \vec{S}_j + J' \sum_{nnn} \vec{S}_i \cdot \vec{S}_j + g\mu_B H \cdot \sum_i \vec{S}_i$$

Application of a magnetic field causes the *Zeeman splitting* of the three excited triplet states  $\rightarrow$  The gap ultimately closes at the critical magnetic field  $H_c = \Delta / (g\mu_B)$



Accordingly magnetic field exceeding  $H_c$  is required to create a finite number of triplets per unit cell of the crystal.

In this depiction:

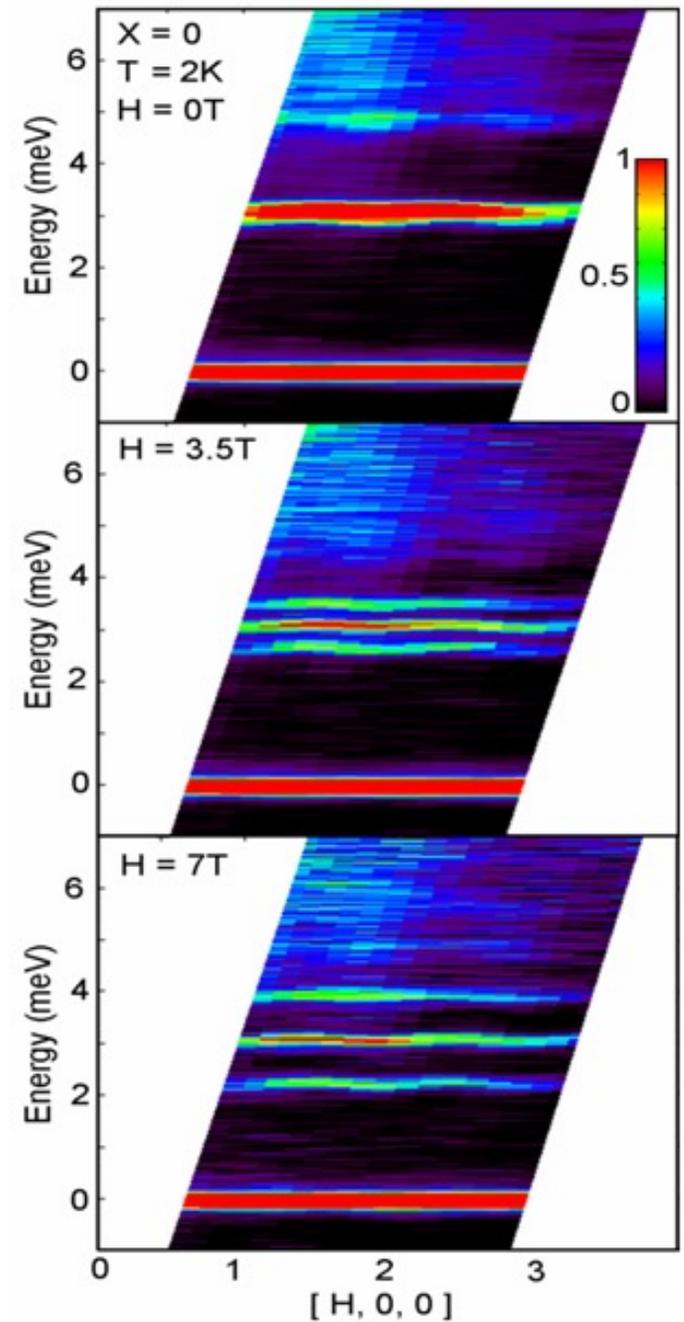
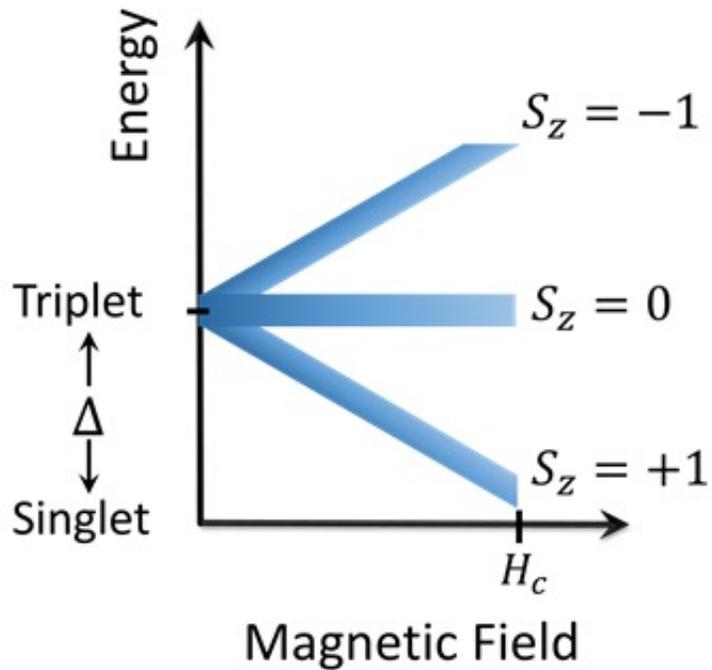
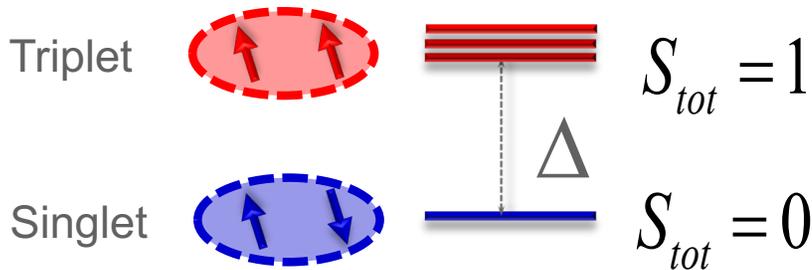
*singlet*  $\leftrightarrow$  *vacuum*

*triplets*  $\leftrightarrow$  *bosons with hard-core repulsion*

*magnetic field*  $\leftrightarrow$  *chemical potential*

# Zeeman Splitting of Triplet States in SCBO

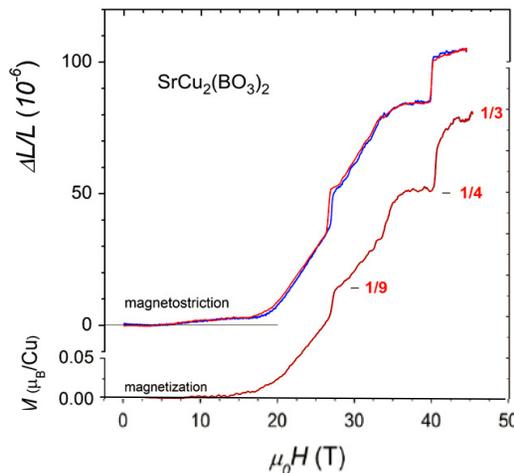
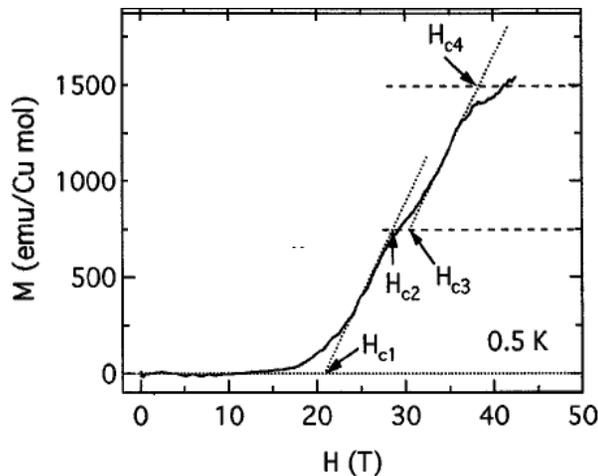
$$\mathcal{H} = J \sum_{nn} \vec{S}_i \cdot \vec{S}_j + J' \sum_{nnn} \vec{S}_i \cdot \vec{S}_j + g\mu_B H \cdot \sum_i \vec{S}_i$$



# Mapping the Phase Diagram > Bound States

## Origin of the magnetization plateaus?

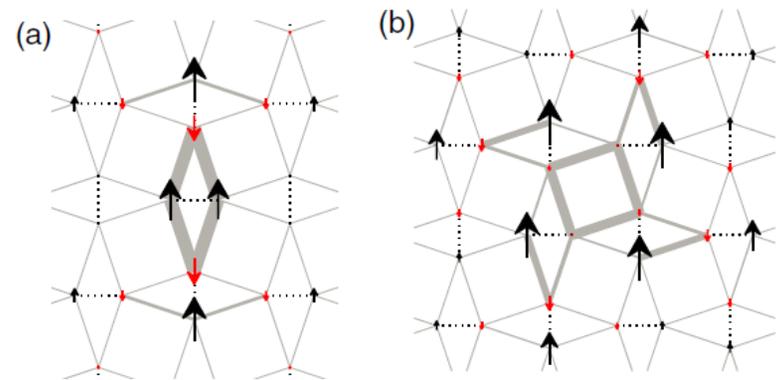
G. Misguich, *et al.*, *Phys. Rev. Lett.* **87**, 097203 (2001)  
 K. Kodama, *et al.*, *Science* **298**, 395 (2002)  
 J. Dorier, *et al.*, *Phys. Rev. Lett.* **101**, 250402 (2008)  
 M. Takigawa, *et al.*, *Phys. Rev. Lett.* **110**, 067210 (2013)  
*and many more...*



H. Kageyama, *et al.*, *Phys. Rev. Lett.* **82**, 3168 (1999)

M. Jaime, *et al.*, *Proc Natl Acad Sci USA* **109**, 12404 (2012)

**Bound states of triplons (pinwheels):  
formation of crystals**



P. Corboz and F. Mila, *Phys. Rev. Lett.* **112**, 147203 (2014)

**Extreme environments needed**

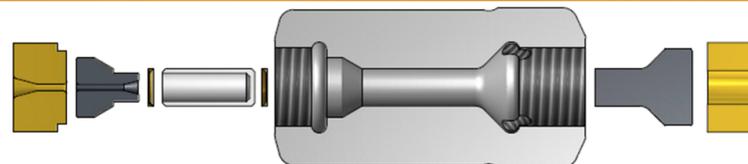
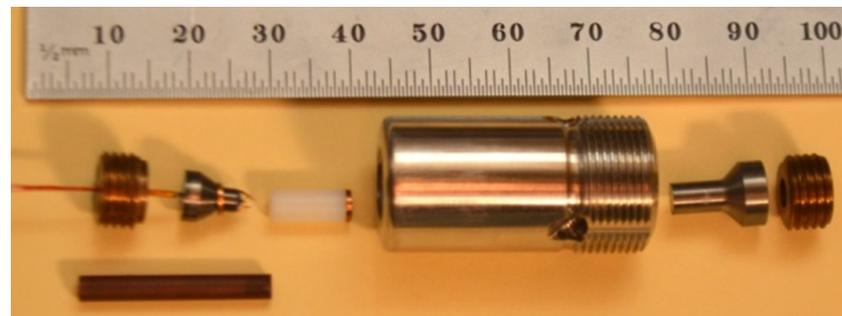
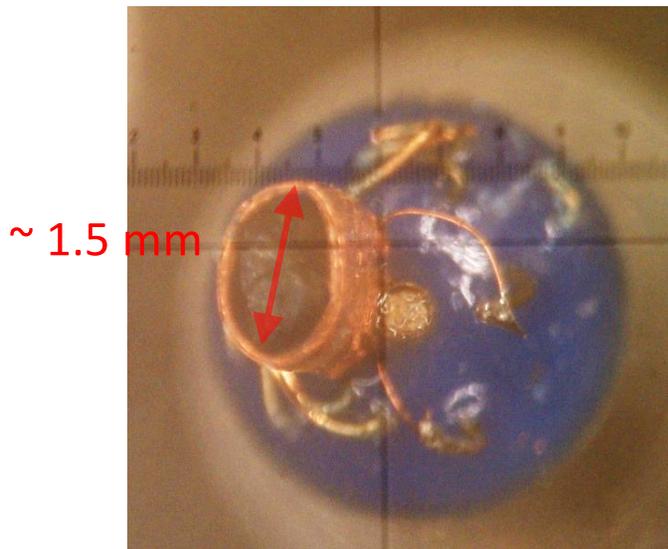
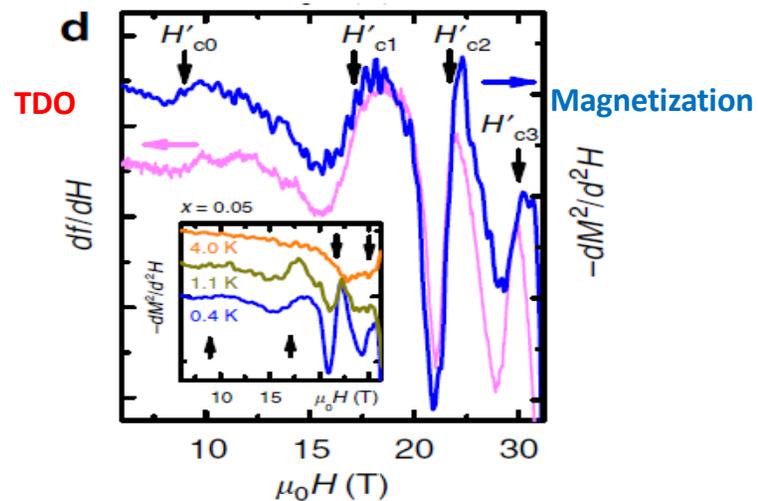
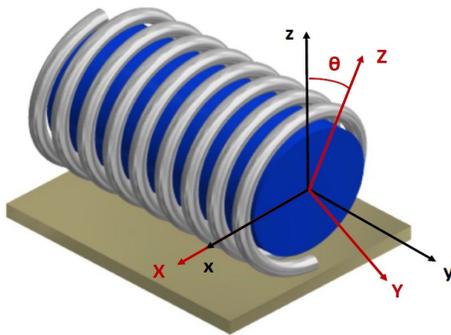


**$T < \sim 2\text{K};$   
 $\mu_0 H \gg \sim 18\text{ T (closing of spin gap)}$**

# Tunnel Diode Oscillator Technique

**Tunnel Diode Oscillators (TDO)** represent a new approach to this challenge. A TDO circuit is constructed of a self-resonating  $LC$  tank biased by a tunnel diode. Sample is located in a coil which acts as the inductor in the  $LC$  tank circuit

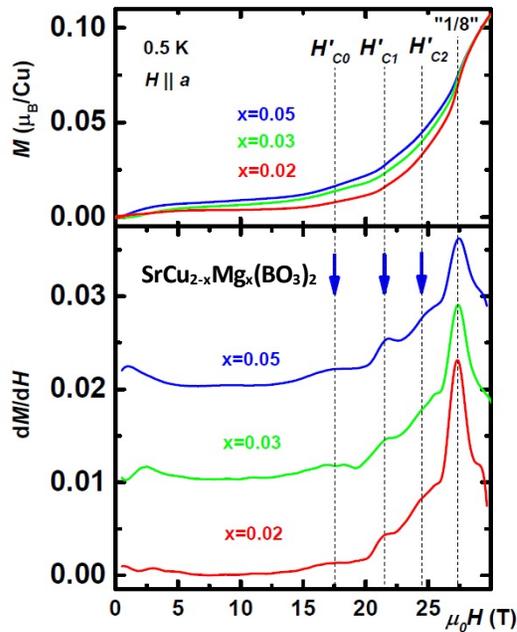
→  $\frac{\Delta f}{f} \propto \frac{dM}{dH}$  enabling us to precisely measure changes in the magnetic moment on the order of  $10^{-12}$  emu



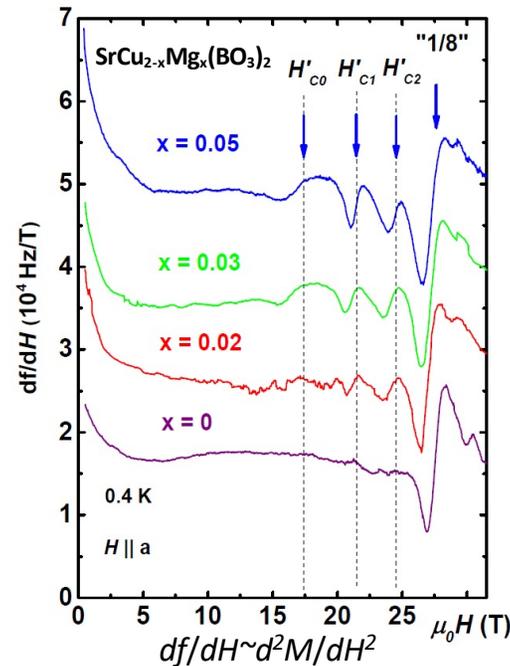
# Tuning knob 3: Chemical Doping

## SrCu<sub>2-x</sub>Mg<sub>x</sub>(BO<sub>3</sub>)<sub>2</sub>: Anomalies below 1/8 plateau

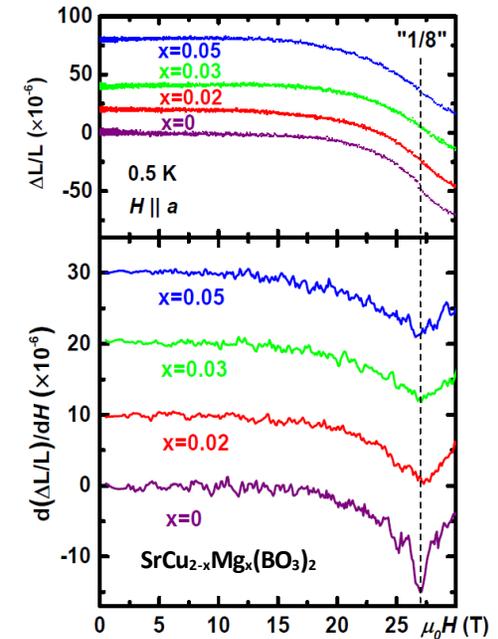
Magnetization measurement



Tunnel diode oscillator (TDO)



Magnetostriction



- Emergent magnetic states at  $H'_{c0} \sim 17.5$  T,  $H'_{c1} \sim 21.5$  T and  $H'_{c2} \sim 24.5$  T
- Absence in magnetostriction: **weak lattice coupling**
- **Superstructure of triplets** with unit cell larger than that for 1/8? **Not likely**

TDO technique:

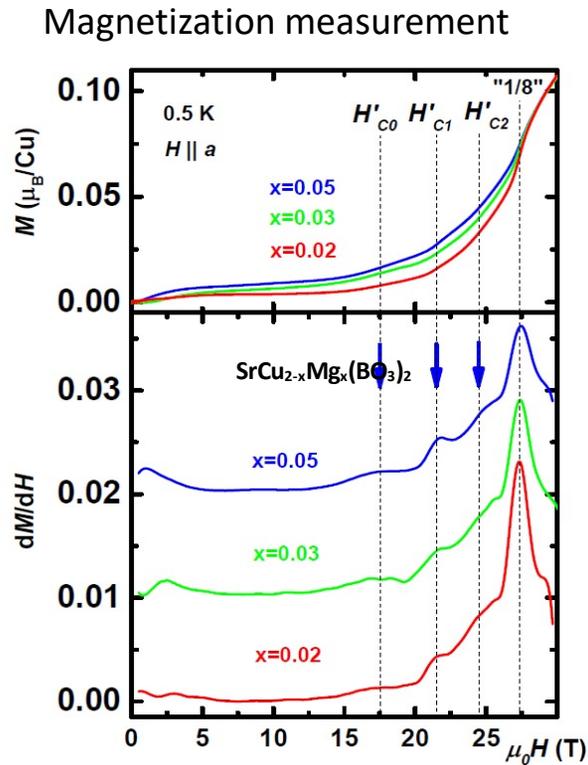
R. Clover, *et al.*, *Rev. Sci. Instrum.* **41**, 617 (1970)

Magnetostriction technique:

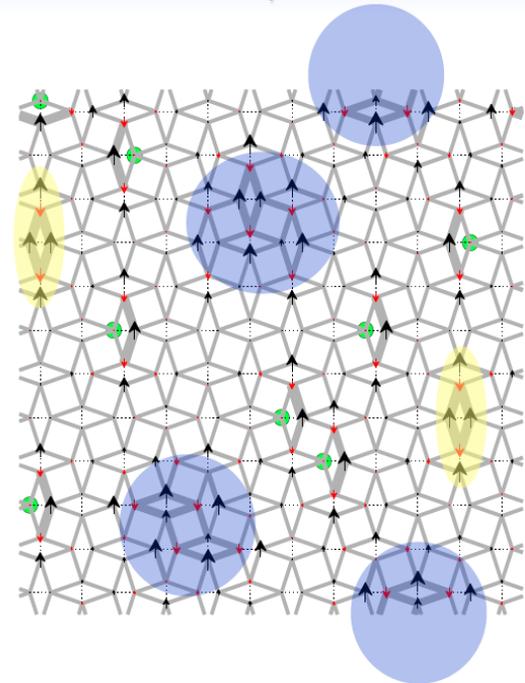
M. Jaime, *Sensor* **17**, 2572 (2017)

Shi, Zhenzhong, *et al.* "Emergent Bound States and Impurity Pairs in Chemically Doped Shastry-Sutherland System." *Nature Commun.* **10**, 2439 (2019).

# Infinite projected entangled-pair state (iPEPS): special configuration of impurities



12x12 with 8 impurities and  $S_z=4+8$



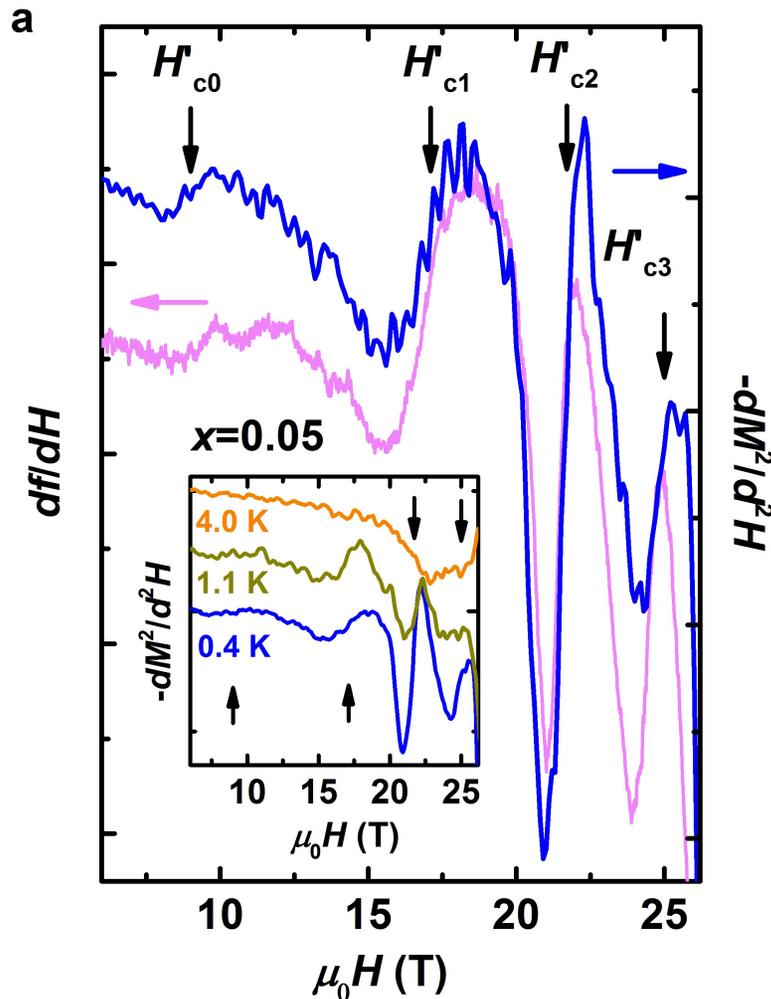
- impurity site  
(inducing  $S_z=1/2$  in the neighborhood)
- bound state ( $S_z=2$ )
- triplet ( $S_z=1$ )

$SrCu_{2-x}Mg_x(BO_3)_2$   
 $x=0.056$

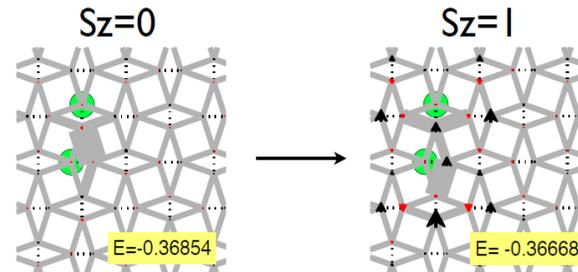
**Better answer: Triplet, Bound states of triplets**

Shi, Zhenzhong, *et al.* "Emergent Bound States and Impurity Pairs in Chemically Doped Shastry-Sutherland System." *Nature Commun.* 10, 2439 (2019).

# Special Configuration of impurities



## Special 2-impurity configurations (8x8 unit cell, D=6)



Energy per excitation:  
**0.2381J**

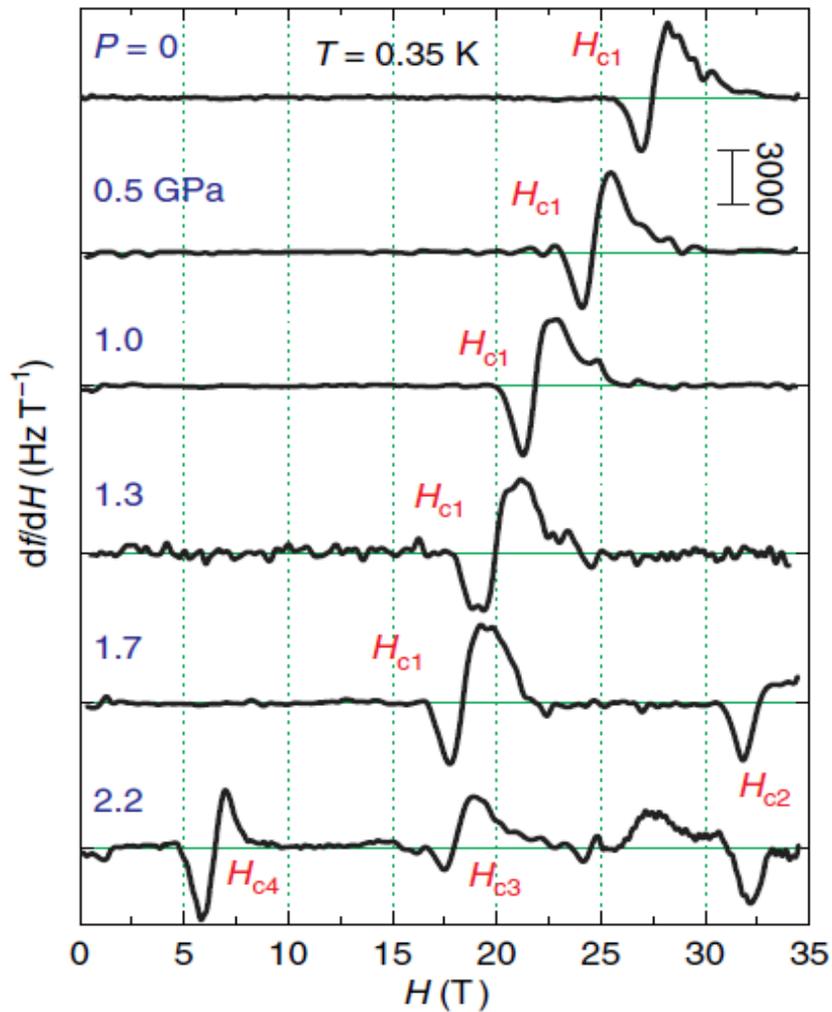
Compared to free triplet excitation:

Energy per excitation:  
**0.430J**

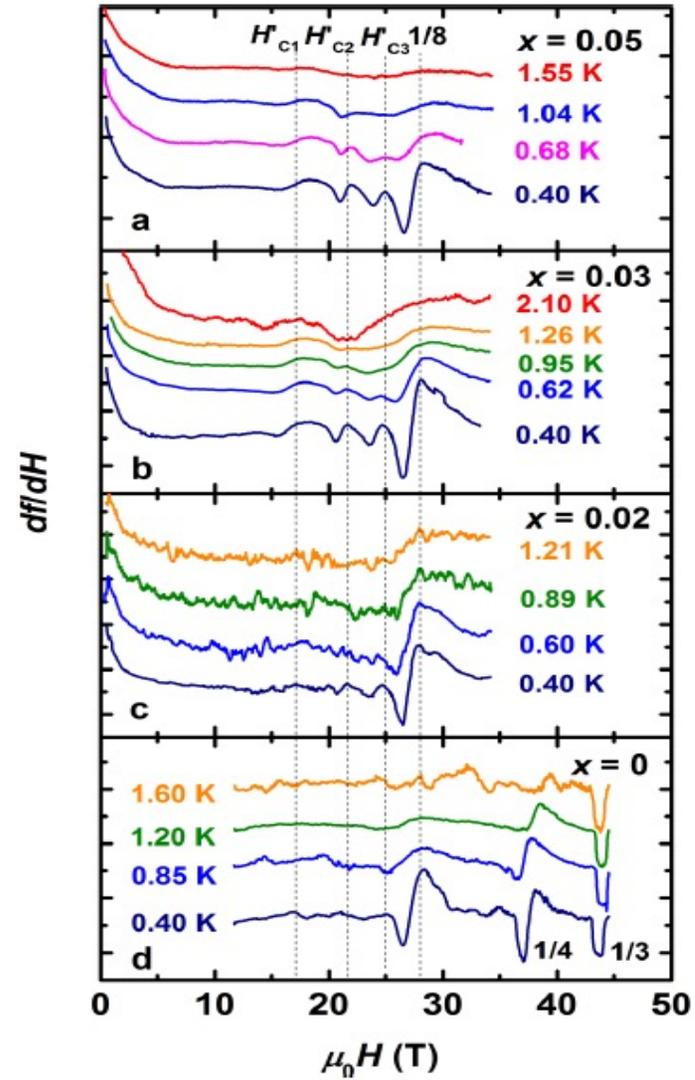
Better answer: **Bound states of impurities**

Shi, Zhenzhong, *et al.* "Emergent Bound States and Impurity Pairs in Chemically Doped Shastry-Sutherland System." *Nature Commun.* 10, 2439 (2019).

# Utilizing all 3 Tuning Knobs



S. Haravifard, *et al.*, *Nat. Commun.* 7, 11956 (2016)



Shi, Zhenzhong, *et al.* *Nature Commun.* 10, 2439 (2019).

# Conclusions

- ❖ Frustrated Magnets are effective candidates in exploring emergent quantum phenomena and topological states of matter.
- ❖ Tuning parameters such as pressure, magnetic field, chemical doping can be used to induce QPTs and map phase diagrams.
- ❖ Technical advancements - specially in sample environment - play key role.